

Computational methods for several models of ice stream flow

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Bathymetry and stickyness distribution

- ▶ Bathymetry:
 - ▶ Aspect ratio $\varepsilon = [H]/[x] \ll 1$
 - ▶ Need surface *and* bed slopes to be small
- ▶ Stickyness distribution:
 - ▶ Limiting cases of plug flow versus vertical shear
 - ▶ Stress ratio: $\lambda = [\tau_{xz}]/[\tau_{\text{membrane}}]$
 - ▶ Discontinuous: frozen to slippery transition at ice stream margins

▶ Bed slope is discontinuous and of order 1.

- ▶ Taylor expansions no longer valid
- ▶ high resolution, subgrid parametrization, short time steps

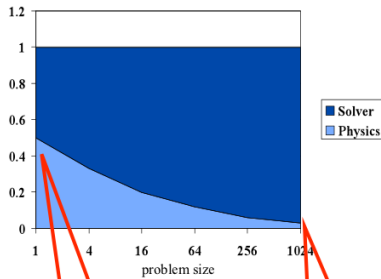
Bathymetry and stickyness distribution

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It's *all* about algorithms (at the petascale)

- **Given, for example:**
 - a “physics” phase that scales as $O(N)$
 - a “solver” phase that scales as $O(N^{3/2})$
 - computation is almost all solver after several doublings
- **Most applications groups have not yet “felt” this curve in their gut**
 - as users actually get into queues with more than 4K processors, this will change

Weak scaling limit, assuming efficiency of 100% in both physics and solver phases



Solver takes 50% time on 128 procs

Solver takes 97% time on 128K procs

(c/o David Keyes)

Hydrostatic equations for ice sheet flow

- ▶ Valid when $w_x \ll u_z$, independent of basal friction (Schoof&Hindmarsh 2010)
- ▶ Eliminate p and w from Stokes by incompressibility:
3D elliptic system for $u = (u, v)$

$$-\nabla \cdot \left[\eta \begin{pmatrix} 4u_x + 2v_y & u_y + v_x & u_z \\ u_y + v_x & 2u_x + 4v_y & v_z \end{pmatrix} \right] + \rho g \bar{\nabla} h = 0$$

$$\eta(\theta, \gamma) = \frac{B(\theta)}{2} (\gamma_0 + \gamma)^{\frac{1-n}{2n}}, \quad n \approx 3$$

$$\gamma = u_x^2 + v_y^2 + u_x v_y + \frac{1}{4} (u_y + v_x)^2 + \frac{1}{4} u_z^2 + \frac{1}{4} v_z^2$$

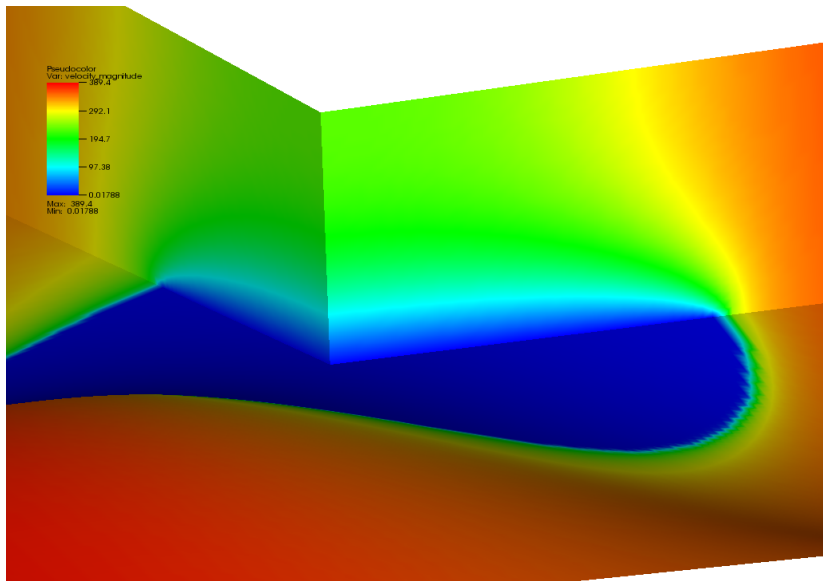
and slip boundary $\sigma \cdot n = \beta^2 u$ where

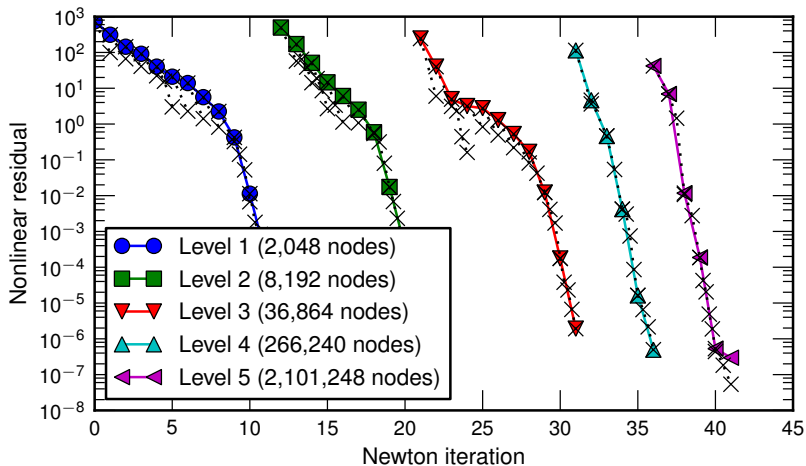
$$\beta^2(\gamma_b) = \beta_0^2 (\varepsilon_b^2 + \gamma_b)^{\frac{m-1}{2}}, \quad 0 < m \leq 1$$

$$\gamma_b = \frac{1}{2} (u^2 + v^2)$$

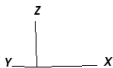
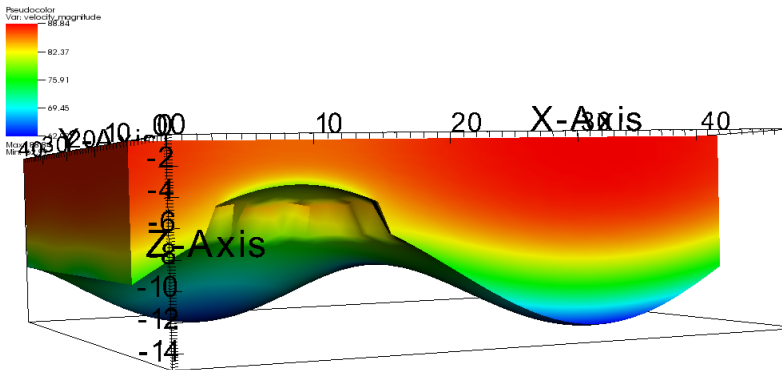
- ▶ Q_1 FEM with Newton-Krylov-Multigrid solver in PETSc:

`src/snes/examples/tutorials/ex48.c`





Grid-sequenced Newton-Krylov solution of test X . The solid lines denote nonlinear iterations, and the dotted lines with \times denote linear residuals.



- ▶ Bathymetry is essentially discontinuous on any grid
- ▶ Shallow ice approximation produces oscillatory solutions
- ▶ Nonlinear and linear solvers have major problems or fail
- ▶ Grid sequenced Newton-Krylov multigrid works as well as in the smooth case

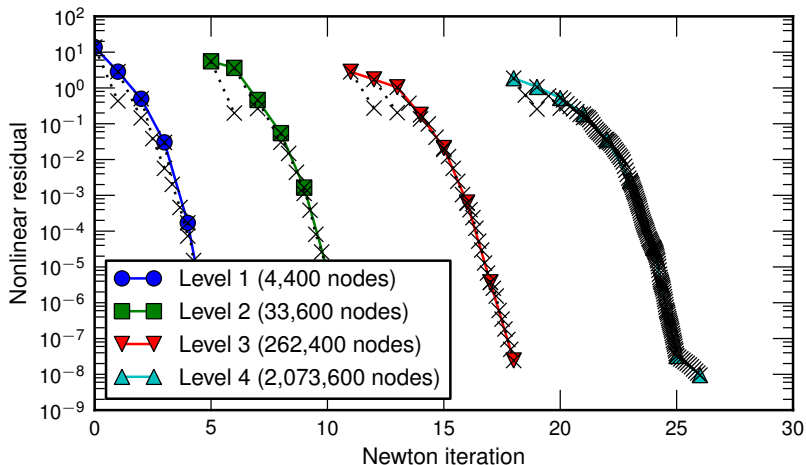


Figure: Grid sequenced Newton-Krylov convergence for test Y . The “cliff” has 58° angle in the red line (12×125 meter elements), 73° for the cyan line (6×62 meter elements).

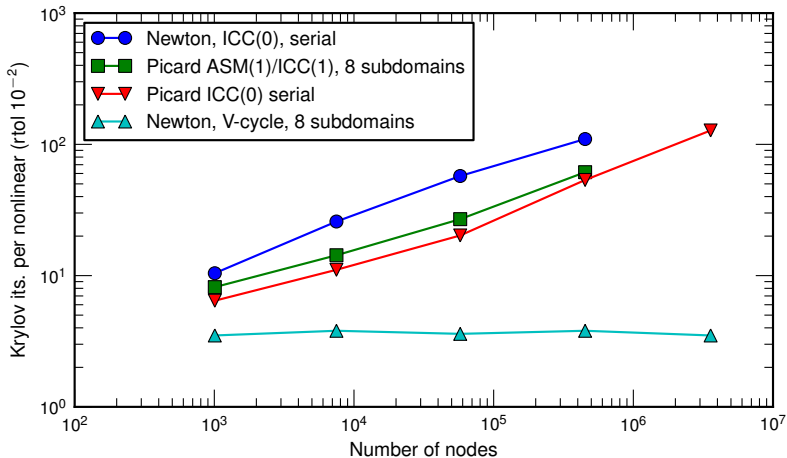
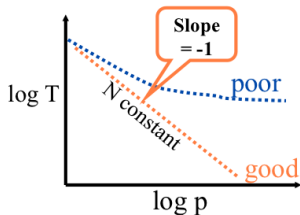


Figure: Average number of Krylov iterations per nonlinear iteration. Each nonlinear system was solved to a relative tolerance of 10^{-2} .

Review: two definitions of scalability

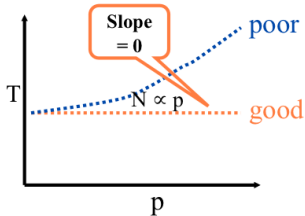
● “Strong scaling”

- execution time (T) decreases in inverse proportion to the number of processors (p)
- *fixed size problem* (N) overall
- often instead graphed as reciprocal, “speedup”



● “Weak scaling” (memory bound)

- execution time remains constant, as problem size and processor number are increased in proportion
- *fixed size problem per processor*
- also known as “Gustafson scaling”



(c/o David Keyes)

Strong scaling on Blue Gene/P (Shaheen)

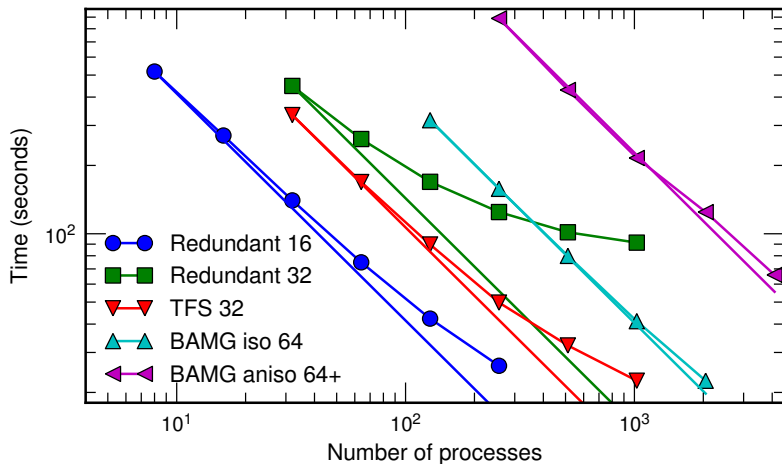


Figure: Strong scaling on Shaheen for different size coarse level problems and different coarse level solvers. The straight lines on the strong scaling plot have slope -1 which is optimal.

Weak scaling on Blue Gene/P (Shaheen)

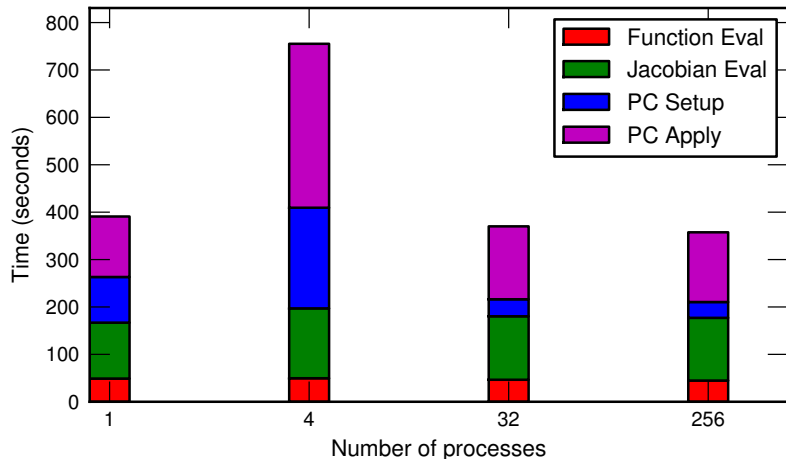


Figure: Weak scaling on Shaheen with a breakdown of time spent in different phases of the solution process. Times are for the full grid-sequenced problem, not just the finest level solve.

One high-accuracy solve costs 30 times as much as a residual evaluation

about 15 to reach truncation error

1000 times faster than existing methods

(Brown, Smith, Ahmadi 2011; submitted to JGR)

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Standard shallow approximations

- ▶ Shallow Ice Approximation (SIA):

- ▶ Purely local definition of velocity

$$u(z) = -(\rho g)^n \int_b^z A(T(z))(h-z)^n \left| \bar{\nabla} h \right|^{n-1} \bar{\nabla} h$$

- ▶ No membrane stresses so only acceptable for $\lambda \gg 1$
- ▶ No solve so costs similar to one residual evaluation per time step

- ▶ “Shelfy Stream” Approximation (SSA)

- ▶ Need to solve elliptic problem posed in the map plane (2D):

$$-\bar{\nabla} \cdot \left[H \bar{\eta} \begin{pmatrix} 4\bar{u}_x + 2\bar{v}_y & \bar{v}_x + \bar{u}_y \\ \bar{v}_x + \bar{u}_y & 2\bar{u}_x + 4\bar{v}_y \end{pmatrix} \right] + \beta^2 \bar{u} + \rho g H \bar{\nabla} h = 0$$

$$\eta(\theta, \gamma) = \frac{B(\theta)}{2} (\gamma_0 + \gamma)^{\frac{1-n}{2n}}, \quad n \approx 3$$

$$\gamma = \bar{u}_x^2 + \bar{v}_y^2 + \bar{u}_x \bar{v}_y + \frac{1}{4} (\bar{u}_y + \bar{v}_x)^2$$

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Vertically-integrated Hybrids

- ▶ Daniel Goldberg 2010, same order of accuracy as hydrostatic
 - ▶ Vertically average “membrane” part of hydrostatic equations

$$-\bar{\nabla} \cdot \left[\bar{\eta} \begin{pmatrix} 4\bar{u}_x + 2\bar{v}_y & \bar{v}_x + \bar{u}_y \\ \bar{v}_x + \bar{u}_y & 2\bar{u}_x + 4\bar{v}_y \end{pmatrix} \right] - \left[\eta \begin{pmatrix} u_z \\ v_z \end{pmatrix} \right]_z + \beta^2 \bar{u} + \rho g H \bar{\nabla} h = 0$$

$$\eta(\theta, \gamma) = \frac{B(\theta)}{2} (\gamma_0 + \gamma)^{\frac{1-n}{2n}}, \quad n \approx 3$$

$$\gamma = \bar{u}_x^2 + \bar{v}_y^2 + \bar{u}_x \bar{v}_y + \frac{1}{4} (\bar{u}_y + \bar{v}_x)^2 + \frac{1}{4} \bar{u}_z^2 + \frac{1}{4} \bar{v}_z^2$$

- ▶ Solve by integrating z -dependence and η , solve linear elliptic problem in map plane for \bar{u} , iterate (Picard, ≈ 50 its)
 - ▶ Evaluating viscosity (or a Newton residual) costs about one hydrostatic residual
- ▶ Bueler and Brown 2009, used in PISM
 - ▶ Ad-hoc combination of independent SSA and SIA solutions
 - ▶ Lower formal order of accuracy, but nonlinear solve is strictly 2D

Non-Newtonian Stokes system: velocity u , pressure p

$$-\nabla \cdot (\eta Du) + \nabla p - f = 0$$

$$\nabla \cdot u = 0$$

$$Du = \frac{1}{2} (\nabla u + (\nabla u)^T)$$

$$\gamma(Du) = \frac{1}{2} Du : Du$$

$$\eta(\gamma) = B(\theta, \dots) (\gamma_0 + \gamma)^{\frac{p-2}{2}}$$

$$p = 1 + \frac{1}{n} \approx \frac{4}{3}$$

$$T = 1 - n \otimes n$$

with boundary conditions

$$(\eta Du - p1) \cdot n = \begin{cases} 0 & \text{free surface} \\ -\rho_w z n & \text{ice-ocean interface} \end{cases}$$

$$u = 0 \quad \text{frozen bed, } \theta < \theta_0$$

$$\left. \begin{aligned} u \cdot n &= g_{\text{melt}}(Tu, \dots) \\ T(\eta Du - p1) \cdot n &= g_{\text{slip}}(Tu, \dots) \end{aligned} \right\} \text{nonlinear slip, } \theta \geq \theta_0$$

$$g_{\text{slip}}(Tu) = \beta_m(\dots) |Tu|^{m-1} Tu$$

Navier $m = 1$, Weertman $m \approx \frac{1}{3}$, Coulomb $m = 0$.

Stokes challenges

Mass conservation is critical

- ▶ Staggered grid finite difference (hard to deal with geometry)
- ▶ Stabilized methods (conservation artifacts when non-smooth)
- ▶ Inf-sup stable mixed finite element method
 - ▶ Use discontinuous pressure to enforce local mass conservation
 - ▶ Inf-sup constant decays like $\sqrt{\varepsilon}$ for $Q_k - P_{k-1}^{\text{disc}}$
 - ▶ Sub-optimal order of accuracy for $Q_k - Q_{k-2}^{\text{disc}}$

Solving saddle-point problems

- ▶ Not uniformly elliptic: solvers are much less robust
- ▶ Standard preconditioners do not work
- ▶ Coupled multigrid with Vanka smoothers offer best performance, not robust for stretched grids or anisotropic viscosity
- ▶ Block preconditioners require approximate commutators, fragile for strong anisotropy and non-smooth viscosity

Outlook

- ▶ We have textbook multigrid efficiency for hydrostatic equations
- ▶ All other models are currently slower at high resolution because there are no scalable implementations
- ▶ Daniel Goldberg's model could be as much as 4 times faster, probably closer to 2 times
- ▶ Bueler & Brown's model (or SSA) could be up to 10 times faster
- ▶ Technical challenges for Stokes
- ▶ Bathymetry is rough enough that we should solve Stokes
- ▶ Singularities: reentrant corners, transition from frozen to slip boundary conditions, grounded margins, grounding lines
- ▶ Implicit time integration

