Tightly coupled solvers with loosely coupled software Modular linear algebra for multi-physics

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Bottlenecks of (Jacobian-free) Newton-Krylov

- Matrix assembly
	- integration/fluxes: FPU
	- insertion: memory/branching
- Preconditioner setup
	- coarse level operators
	- overlapping subdomains
	- (incomplete) factorization
- Preconditioner application
	- triangular solves/relaxation: memory
	- coarse levels: network latency
- Matrix multiplication
	- Sparse storage: memory
	- Matrix-free: FPU

Globalization

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Hardware capabilities

Floating point unit

Recent Intel: each core can issue

- 1 packed add (latency 3)
- 1 packed mult (latency 5)
- One can include an aligned read
- Out of Order execution
- Peak: 10 Gflop/s (double)

Memory

- ∼ 250 cycle latency
- 5.3 GB/s bandwidth
- 1 double load / 3.7 cycles
- Pay by the cache line (32/64 B)
- • L2 cache: \sim 10 cycle latency

Memory Bandwidth

• Stream Triad benchmark (GB/s): $w \leftarrow \alpha x + y$

Threads per Node	Cray XT5		BlueGene/P	
	Total	Per Core	Total	Per Core
	8448	8448	2266	2266
2	10112	5056	4529	2264
4	10715	2679	8903	2226
6	10482	1747	$\overline{}$	\sim

• Sparse matrix-vector product: 6 bytes per flop

Sparse Mat-Vec performance model

Compressed Sparse Row format (AIJ)

For $m \times n$ matrix with N nonzeros

- ai row starts, length $m+1$
- aj column indices, length *N*, range [0,*n*−1)

aa nonzero entries, length *N*, scalar values

$$
y \leftarrow y + Ax
$$

\nfor (i=0; i \leftarrow n; i++)
\nfor (j=ai[i]; j\n $y[i] += aa[j] * x[a[j]]$;

- One add and one multiply per inner loop
- Scalar aa [j] and integer $a \in [j]$ only used once
- Must load $a \dagger$ [\dagger] to read from x, may not reuse cache well

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Optimizing Sparse Mat-Vec

- Order unknows so that vector reuses cache (Reverse Cuthill-McKee)
	- Optimal: $\frac{(2 \text{ flops})(\text{bandwidth})}{\text{sizeof}(\text{Scalar}) + \text{sizeof}(\text{Int})}$
	- Usually improves strength of ILU and SOR
- Coalesce indices for adjacent rows with same nonzero pattern (Inodes)
	- Optimal: (2 flops)(bandwidth) sizeof(Scalar)+sizeof(Int)/*i*
	- Can do block SOR (much stronger than scalar SOR)
	- Default in PETSc, turn off with -mat_no_inode
	- Requires ordering unknowns so that fields are interlaced, this is (much) better for memory use anyway
- Use explicit blocking, hold one index per block (BAIJ format)
	- Optimal: $\frac{(2 \text{ flops})(\text{bandwidth})}{\text{sizeof}(\text{Scalar}) + \text{sizeof}(\text{Int})/b^2}$
	- Block SOR and factorization
	- Symbolic factorization works with blocks (much cheaper)
	- Very regular memory access, unrolled dense kernels
	- Faster insertion: MatSetValuesBlocked()

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Performance of blocked matrix formats

Throughput (Mflop/s) for different matrix formats on Core 2 Duo (P8700) and Opteron 2356 (two sockets). MatSolve is a forward- and back-solve with incomplete Cholesky factors. The AIJ format is using "inodes" which unrolls across consecutive rows with identical nonzero pattern (pairs in this case).

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Optimizing unassembled Mat-Vec

• High order spatial discretizations do more work per node

- Dense tensor product kernel (like small BLAS3)
- Cubic (Q_3) elements in 3D can achieve $> 70\%$ of peak FPU (compare to $< 6\%$ for assembled operators on multicore)
- Can store Jacobian information at quadrature points (usually pays off for Q_2 and higher in 3D)
- Spectral, WENO, DG, FD
- Often still need an assembled operator for preconditioning
- Boundary element methods
	- Dense kernels
	- Fast Multipole Method (FMM)

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Power-law Stokes Scaling

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What you can do

- Speak at the most specific language possible
	- 3D structural analysis: symmetric block size 3
	- 3D compressible flow: nonsymmetric block size 5
- Order unknowns for cache reuse (low-bandwidth like RCM is good)
- Dual order
	- Assemble a low-order discretization
	- Provide matrix-free high-order operator (FD, ADI, caching at quadrature points)
	- More robust with SOR and ILU due to *h*-ellipticity
	- Sometimes Picard linearization has a more compact stencil

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Moussau et al, 2002.

Linear reaction-diffusion, split method converges to the wrong steady state . Knoll et al, 2003.

- CFL too restrictive for explicit
	- But hyperbolic systems do not weak scale if you care about phase
- Naive semi-implicit has poor accuracy, stability, robustness
- Good IMEX exists, but still need to treat stiff part implicitly

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Coupled approach to multiphysics

- Smooth all components together
	- Block SOR is the most popular
	- Vanka smoothers for indefinite problems
	- Block II U often more robust
- Scaling between fields is critical
- **Indefiniteness**
	- Make smoothers and interpolants respect inf-sup condition
	- Difficult to handle anisotropy
	- Can use Schur field-split to define a smoother
- Transport
	- Define smoother in terms of first-order upwind discretization (*h*-ellipticity)
	- Evaluate residuals using high-order discretization
	- Use Schur field-split to "parabolize" at the top level or to define smoother on levels
- • Open research area, hard to write modular software

Anisotropy, Heterogeneity

• Anisotropy

- Semi-coarsening
- Line smoothers
- Order unknowns so that incomplete factorization "includes" a line smoother

• Heterogeneity

- Make coarse grids align
- Strong smoothers
- Energy-minimizing interpolants
- Mostly possible with generic software components

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Splitting for Multiphysics

$$
\begin{bmatrix} A & B \\ C & D \end{bmatrix} \begin{bmatrix} x \\ y \end{bmatrix} = \begin{bmatrix} f \\ g \end{bmatrix}
$$

Relaxation: -pc fieldsplit type [additive,multiplicative, symmetric_multiplicative]

$$
\begin{bmatrix} A & D \end{bmatrix}^{-1} \qquad \begin{bmatrix} A & D \end{bmatrix}^{-1} \qquad \begin{bmatrix} A & D \end{bmatrix}^{-1} \begin{bmatrix} 1 - \begin{bmatrix} A & B \end{bmatrix} \begin{bmatrix} A & D \end{bmatrix}^{-1} \end{bmatrix}
$$

- Gauss-Seidel inspired, works when fields are loosely coupled
- Factorization: -pc fieldsplit type schur

$$
\begin{bmatrix} A & B \\ S \end{bmatrix}^{-1} \begin{bmatrix} 1 & 1 \\ CA^{-1} & 1 \end{bmatrix}^{-1}, \qquad S = D - CA^{-1}B
$$

- robust (exact factorization), can often drop lower block
- how to precondition *S* which is usually dense?
	- interpret as differential operators, use appro[xim](#page-15-0)[ate](#page-17-0) [c](#page-15-0)[om](#page-16-0)[m](#page-11-0)[ut](#page-12-0)[a](#page-23-0)[to](#page-24-0)[r](#page-11-0)[s](#page-12-0)

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Physics-based preconditioners (semi-implicit method)

Shallow water with stiff gravity wave

^h is hydrostatic pressure, *^u* is velocity, [√] *gh* is fast wave speed

$$
h_t - (uh)_x = 0
$$

$$
(uh)_t + (u^2h + \frac{1}{2}gh^2)_x = 0
$$

Semi-implicit method

Suppress spatial discretization, discretize in time, implicitly for the terms contributing to the gravity wave

$$
\frac{h^{n+1} - h^n}{\Delta t} + (uh)_x^{n+1} = 0
$$

$$
\frac{(uh)^{n+1} - (uh)^n}{\Delta t} + (u^2h)_x^n + g(h^n h^{n+1})_x = 0
$$

Rearrange, eliminating (*uh*) *n*+1

$$
\frac{h^{n+1}-h^n}{\Delta t}-\Delta t(gh^nh_x^{n+1})_x=-S_x^n
$$

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Delta form

• Preconditioner should work like the Newton step

$$
-F(x) \mapsto \delta x
$$

• Recast semi-implicit method in delta form

$$
\frac{\delta h}{\Delta t} + (\delta u h)_x = -F_0, \quad \frac{\delta u h}{\Delta t} + gh^n(\delta h)_x = -F_1, \quad \hat{J}\begin{pmatrix} \frac{1}{\Delta t} & \text{div} \\ gh^n \nabla & \frac{1}{\Delta t} \end{pmatrix}
$$

• Eliminate δ*uh*

$$
\frac{\delta h}{\Delta t}-\Delta t(gh^n(\delta h)_x)_x=-F_0+(\Delta t F_1)_x,\quad S\sim \frac{1}{\Delta t}-g\Delta t\,\text{div}\,h^n\nabla
$$

• Solve for δ*h*, then evaluate

$$
\delta u h = -\Delta t \big[gh^n(\delta h)_x - F_1 \big]
$$

- Fully implicit solver
	- Is nonlinearly consistent (no splitting error), can be high-order in time
	- Leverages existing code when a semi-implicit method has been implemented
	- Allows efficient bifurcation analysis, steady-state analysis

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Stokes

Weak form of the Newton step

Find (u, p) such that

$$
\int_{\Omega} (Dv)^{T} [\eta 1 + \eta' Dw \otimes Dw] Du
$$

- $p \nabla \cdot v - q \nabla \cdot u = -v \cdot F(w)$ $\forall (v,q)$

Matrix

$$
\begin{bmatrix} A(w) & B^T \ B & \end{bmatrix} \begin{pmatrix} u \ p \end{pmatrix} = - \begin{pmatrix} F_u(w) \ 0 \end{pmatrix}
$$

Block factorization

$$
\begin{bmatrix} A & B^T \ B & \end{bmatrix} = \begin{bmatrix} 1 & 1 \ BA^{-1} & 1 \end{bmatrix} \begin{bmatrix} A & B^T \ B & S \end{bmatrix} = \begin{bmatrix} A & 1 \ BA & S \end{bmatrix} \begin{bmatrix} 1 & A^{-1}B^T \ 1 & 1 \end{bmatrix}
$$

where the Schur complement is

$$
S=-BA^{-1}B^T.
$$

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Properties of the Schur complement

Block factorization

where

$$
\begin{bmatrix} A & B^T \ B & A \end{bmatrix} = \begin{bmatrix} 1 & 1 \ BA^{-1} & 1 \end{bmatrix} \begin{bmatrix} A & B^T \ S & B \end{bmatrix} = \begin{bmatrix} A & 1 \ BA & 1 \end{bmatrix}
$$

$$
S = -BA^{-1}B^T.
$$

- *S* is symmetric negative definite if *A* is SPD and *B* has full rank (discrete inf-sup condition)
- *S* is dense
- We only need to multiply B, B^T with vectors.
- We need preconditioners for *A* and *S*.
- Any definite preconditioner can be used for *A*.
- It's not obvious how to precondition *S*, more on that later.

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Preconitioning the Schur complement

• $S = -BA^{-1}B^{T}$ is dense so we can't form it, we need S^{-1} .

Physics-based commutator: anisotropic pressure diffusion

$$
v^T A(w)u \sim \int (Dv)^T [\eta 1 + \eta' Dw \otimes Dw] Du
$$

• We would like to find an operator A_p such that

$$
-S = BA^{-1}B^T \approx BB^TA_p^{-1} =: P_S
$$

so that

$$
P_S^{-1}=A_p(BB^T)^{-1}
$$

• Note

$$
\textit{BB}^\mathcal{T} \sim (-\nabla \cdot) \nabla = -\Delta
$$

corresponds to a Laplacian in the pressure space (multigrid).

 \bullet If $\eta', \nabla \eta \ll 1$ then $A_\rho \sim -\eta \Delta$ so $P_S^{-1} = \eta 1$

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Least squares commutator

• Schur complement

$$
S=-BA^{-1}B^T
$$

Suppose *B* is square and nonsingular. Then

$$
S^{-1}=-B^{-T}AB^{-1}.
$$

B is not square, replace B^{-1} with Moore-Penrose pseudoinverse

$$
B^{\dagger} = B^{T} (BB^{T})^{-1}
$$
, $(B^{T})^{\dagger} = (BB^{T})^{-1}B$.

Then

$$
P_S^{-1} = -(BB^T)^{-1}BAB^T(BB^T)^{-1}.
$$

- Requires 2 Poisson preconditioners for $(BB^T)⁻¹$ per iteration
- Better with scaling, from mass matrices and effective viscosity (Elman et al. 2006, May & Moresi 2008)
- -pc_type fieldsplit -pc_fieldsplit_type schur -fieldsplit p_pc_type_lsc -fieldsplit_p_lsc_pc_type mg AD > x E > x E > +

Unsteady Navier-Stokes

Strong form

$$
J(w)\begin{bmatrix}u\\p\end{bmatrix}\sim\begin{cases}\rho(\alpha u+w\cdot\nabla u+u\cdot\nabla w)-\eta\nabla^2 u+\nabla p=-F(w)\\ \nabla\cdot u=0\end{cases}
$$

Matrix form

$$
\begin{bmatrix} A(w) & B^T \ B & \end{bmatrix} = \begin{bmatrix} 1 & 1 \ BA^{-1} & 1 \end{bmatrix} \begin{bmatrix} A & B^T \ B & \end{bmatrix} \qquad S = -BA^{-1}B^T
$$

Define *A*(*w*) in pressure space

- Want $P_S = (BB^T)A_p^{-1} \approx BA^{-1}B^T$, $P_S^{-1} = A_p(BB^T)^{-1}$
- A_ρ ∼ρ $\left(\alpha p + w\cdot \nabla \rho + \rho \text{tr}(\nabla w)\right) \eta \nabla^2 \rho$
- *p* tr(∇*w*) term is questionable, not needed for Picard
- • Almost mesh-independent, weak Reynolds number dependence

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Flow Control for a PETSc Application

Overwhelmed with choices

- If you have a hard problem, no black-box solver will work well
- Everything in PETSc has a plugin architecture
	- Put in the "special sauce" for your problem
	- Your implementations are first-class
- PETSc exposes an algebra of composition at runtime
	- Build a good solver from existing components, at runtime
	- Multigrid, domain decomposition, factorization, relaxation, field-split
	- Choose matrix format that works best with your preconditioner
	- structural blocking, Neumann matrices, monolithic versus nested

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- package each "physics" independently
- solve single-physics and coupled problems
- semi-implicit and fully implicit
- reuse residual and Jacobian evaluation unmodified
- direct solvers and efficient fieldsplit without recompilation
- use the best possible matrix format for each physics (e.g. symmetric block size 3)
- matrix-free anywhere
- multip[le](#page-26-0) l[ev](#page-28-0)[e](#page-26-0)[ls](#page-27-0)[o](#page-32-0)[f](#page-23-0) [n](#page-24-0)[es](#page-34-0)[ti](#page-23-0)n[g](#page-34-0)

Velocity Stokes Pressure

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Boundary Layer

Ocean

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MatNest: a matrix format for field-split

- pseudo-elasticity for mesh motion
- $(\dot{x} u) \cdot n =$ accumulution
- "just" geometry
- Stokes problem
- temperature dependence of rheology
- ALE and strain heating in heat transport
- thermal advection-diffusion
- Blocks stored separately no-copy access
- MatGetSubMatrix API
	- looks same as "normal" matrices
- Nesting can be recursive

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• Implements standard linear algebra operations

MatGetLocalSubMatrix(Mat A,IS rows,IS cols,Mat *B);

- Primarily for assembly
	- B is not guaranteed to implement MatMult
	- \bullet The communicator for B is not specified, only safe to use non-collective ops (unless you check)
- IS represents an index set, includes a block size and communicator
- MatSetValuesBlockedLocal() is implemented
- MatNest returns nested submatrix, no-copy
- No-copy for Neumann-Neumann formats (unassembled across procs, e.g. BDDC, FETI-DP)
- Most other matrices return a lightweight proxy Mat
	- COMM_SELF
	- Values not copied, does not implement MatMult
	- Translates indices to the language of the parent matrix
	- Multiple levels of nesting are flattened

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Wrap-up

- Software modularity while retaining access to good solvers
	- Reuse single-physics modules
	- Unintrusive "special sauce" (once you figure it out)
- Choose the matrix format at runtime, best for your preconditioner
	- monolithic, nested, Neumann
	- scalar or block, symmetric
- • Break into pieces that are "understood", keep some block structure for high throughput