A parallel unstructured implicit 3D polythermal model for outlet glaciers

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Conservative models

Conforming boundaries and model non-smoothness

Multilevel Solvers



Bathymetry and stickyness distribution

- Bathymetry:
 - Aspect ratio $\varepsilon = [H]/[x] \ll 1$
 - Need surface and bed slopes to be small
- Stickyness distribution:
 - Limiting cases of plug flow versus vertical shear
 - Stress ratio: $\lambda = [\tau_{xz}]/[\tau_{membrane}]$
 - Discontinuous: frozen to slippery transition at ice stream margins
- Standard approach in glaciology: Taylor expand in ε and sometimes λ , drop higher order terms.
- $\lambda \gg 1$ Shallow Ice Approximation (SIA), no horizontal coupling
- $\lambda \ll 1~$ Shallow Shelf Approximation (SSA), 2D elliptic solve in map-plane
 - Hydrostatic and various hybrids, 2D or 3D elliptic solves
- Bed slope is discontinuous and of order 1.
 - Taylor expansions no longer valid
 - Numerics require high resolution, subgrid parametrization, short time steps
 - Still get low quality results in the regions of most interest.
- Basal sliding parameters are discontinuous.



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Polythermal ice

- Interface tracking methods (e.g. Greve's SICOPOLIS)
 - Different fields for temperate and cold ice.
 - Lagrangian or Eulerian, problems with changing topology
 - No discrete conservation
- Interface capturing
 - Enthalpy: Aschwanden, Bueler, Khroulev, Blatter (J. Glac. 2012/PISM)
 - Not in conservation form
 - Only conservative for infinitesimal melt fraction
 - Energy
 - Conserves mass, momentum, and energy for arbitrary melt fraction

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Implicit equation of state

Conservative (non-Boussinesq) two-phase ice flow

Find momentum density ρu , pressure p, and total energy density E:

$$(\rho u)_t + \nabla \cdot (\rho u \otimes u - \eta D u_i + p 1) - \rho g = 0$$

$$\rho_t + \nabla \cdot \rho u = 0$$

$$E_t + \nabla \cdot ((E+p)u - k_T \nabla T - k_\omega \nabla \omega) - \eta D u_i : D u_i - \rho u \cdot g = 0$$

- Solve for density ρ, ice velocity u_i, temperature T, and melt fraction ω using constitutive relations.
 - Simplified constitutive relations can be solved explicitly.
 - Temperature, moisture, and strain-rate dependent rheology η.
 - ► High order FEM, typically Q₃ momentum & energy
- DAEs solved implicitly after semidiscretizing in space.
- Preconditioning using nested fieldsplit
- Thermomechanical steady state in about 10 nonlinear iterations



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Why care about conforming and non-smoothness?

- Subshelf ocean circulation is physically richer and more difficult to solve than ice flow
 - $\blacktriangleright\,$ Heat transfer is sensitive to boundary layer processes with thickness < 1 m, $Re \sim 10^6$
 - Countless engineering studies: wall modeling is limited, significant normal resolution still necessary
 - Unaligned interface anisotropy is bad for Eulerian AMR methods
- We care about high-dimensional sensitivity analysis and inversion
 - Forward model evaluations alone are an inefficient way to explore a high-dimensional space
 - Each source of model non-smoothness requires a lot of analysis to use adjoint methods
 - Conforming moving mesh methods eliminate all but "essential" non-smoothness (like contact)

Mesh motion via Inverse Beltrami formulation

- nonlinear elliptic (or parabolic) equation for mesh location
- prescribe resolution and anisotropy using target metric tensor
- efficient solution using Newton-Krylov multigrid or nonlinear multigrid

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- conservative slip boundary conditions at most surfaces
- ALE transport scheme corrected to satisfy geometric conservation law

Transport

Theorem (Godunov 1954)

Non-oscillatory linear spatial discretizations for transport are at most first order accurate.

- First order accurate discretizations have unacceptably high numerical diffusion
- Discretization choices:
 - Limited or reconstructed finite volume or finite difference
 - Second order TVD limiters have corners
 - Weighted Essential Non-Oscillatory (WENO) smooth, but very nonlinear

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- Central or upwind
- Inconvenient for unstructured grids
- Nonlinearly stabilized continuous finite element
 - Currently used, but fragile and messy
- Discontinuous Galerkin
 - Cell-wise entropy stability without limiters
 - Improved robustness with smooth limiters
 - Transitioning to this

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Multilevel Solvers

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Why do we need multilevel solvers?

- Elliptic problems are globally coupled
- Without a coarse level, number of iterations proportional to inverse mesh size
- High-volume local communication is an inefficient way to communicate long-range information, bad for parallel models
- Most important with 3D flow features and/or slippery beds
- Nested/split multilevel methods
 - Decompose problem into simpler sub-problems, use multilevel methods on each
 - Good reuse of existing software
 - More synchronization due to nesting, more suitable after linearization
- Monolithic/coupled multilevel methods
 - Better convergence and lower synchronization, but harder to get right
 - Internal nonlinearities resolved locally
 - ► More discretization-specific, less software reuse

Relative effect of the blocks

$$J = \begin{pmatrix} J_{uu} & J_{up} & J_{uE} \\ J_{pu} & 0 & 0 \\ J_{Eu} & J_{Ep} & J_{EE} \end{pmatrix}.$$

- *J_{uu}* Viscous/momentum terms, nearly symmetric, variable coefficients, anisotropy from Newton.
- J_{up} Weak pressure gradient, viscosity dependence on pressure (small), gravitational contribution (pressure-induced density variation). Large, nearly balanced by gravitational forcing.
- J_{uE} Viscous dependence on energy, very nonlinear, not very large.
- J_{pu} Divergence (mass conservation), nearly equal to J_{up}^{T} .
- J_{Eu} Sensitivity of energy on momentum, mostly advective transport. Large in boundary layers with large thermal/moisture gradients.
- J_{Ep} Thermal/moisture diffusion due to pressure-melting, $u \cdot \nabla$.
- *J_{EE}* Advection-diffusion for energy, very nonlinear at small regularization. Advection-dominated except in boundary layers and stagnant ice, often balanced in vertical.

How much nesting?

$$P_{1} = \begin{pmatrix} J_{uu} & J_{up} & J_{uE} \\ 0 & B_{pp} & 0 \\ 0 & 0 & J_{EE} \end{pmatrix}$$

- *B_{pp}* is a mass matrix in the pressure space weighted by inverse of kinematic viscosity.
- Elman, Mihajlović, Wathen, JCP 2011 for non-dimensional isoviscous Boussinesq.
- Works well for non-dimensional problems on the cube, not for realistic parameters.

$$P = \begin{bmatrix} \begin{pmatrix} J_{uu} & J_{up} \\ J_{pu} & 0 \end{pmatrix} & \\ \begin{pmatrix} J_{Eu} & J_{Ep} \end{pmatrix} & J_{EE} \end{bmatrix}$$

- Inexact inner solve using upper-triangular with B_{pp} for Schur.
- Another level of nesting.
- GCR tolerant of inexact inner solves.
- Outer converges in 1 or 2 iterations.
- Low-order preconditioning full-accuracy unassembled high order operator.
- Build these on command line with PETSc PCFieldSplit.

Full Approximation Scheme

$$\begin{split} \tilde{u}^{h} &\leftarrow S^{h}_{\mathsf{pre}} u^{h}_{0} \\ L^{H} u^{H} &= I^{H}_{h} f^{h} + \underbrace{L^{H} \hat{I}^{H}_{h} \tilde{u}^{h} - I^{H}_{h} L^{h} \tilde{u}^{h}}_{\tau^{H}_{h}} \\ u^{h} &\leftarrow S^{h}_{\mathsf{post}} \Big[\tilde{u}^{h} + I^{h}_{H} (u^{H} - \hat{I}^{H}_{h} \tilde{u}^{h}) \Big] \end{split}$$

pre-smooth solve coarse problem for u^H

apply correction and post-smooth

- Nonlinearities from spatial discretization fixed locally
- No assembled matrices so better floating point utilization, less memory
- Makes progress on all physical components at once
- FD and DG good, less efficient for continuous finite element methods
- Influence of surface evolution is low rank, no need to visit finest level on each iteration

Outlook

- Basal hydrology model
 - Need mesh-independent statistics
 - Numerical homogenization?
- True inverse and sensitivity support, but need to invert for the right thing
- Dynamic remeshing after large topology changes
- Finish FAS multigrid for full coupled system including geometry
- User-friendliness of process
 - Currently using georeferenced initial/boundary data via GDAL

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- But meshing process is not fully automatic
- Better: multiresolution database (Mark Fahnestock)